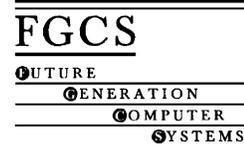




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## Kites flying in and out of space—distributed physically based art on the grid

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### Abstract

In this paper, we describe the design and implementation of a Virtual Reality (VR) art piece—“Kites flying in and out of space” that was inspired by the kite-like art forms of French artist, Jackie Matisse. We use a physically based animation method known as the mass-spring model to realistically simulate the movement of these virtual kitetail forms in the CAVE VR theatre. In this immersive environment, the user can interact with these “virtual” kites by moving them, changing their imagery or adding a wind force. However, the real-time requirements imposed by immersive environments and the computational complexity in calculating these forms inhibit the number of kites we can “fly”. To address this limitation, we show how the use of distributed computing resources across the GRID can provide a scalable solution. Serendipitously, we also discovered that the movement of these virtual art forms became visual metaphors for the network performance and parameters.

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*Keywords:* Virtual Reality; GRID Art; Kitetails

### 1. Introduction

#### 1.1. Motivation—Jacqueline Matisse’ kitetails

French sculpter and light artist, Jackie Matisse [1] creates Teflon or crepe kites, with artistic tails as long as 15 ft, that can soar through the air, ripple through water, or undulate with the air currents in a room. Randomly influenced by natural forces, the kitetails move, and metamorphose in faint air currents and dramatically changing natural light; echoing the more intense pressures of “civilized” life, they interact with visitors who traverse the gallery.

The VR piece was inspired by the three-screen collaborative video *Sea Tails* created in 1983 by Matisse with filmmaker Molly Davies. Fig. 1 shows a still from this video. The film follows 10 kitetails on their dancing flight through the air and into the water. Our goal here was to realistically simulate the movement of these physical kite forms in a virtual environment.

#### 1.2. Background work

To realistically simulate the movement of these kitetails, we draw upon existing research in physically based cloth animation [2]. Animation in an immersive environment needs to be robust and fast given the high-frame rate and interaction requirements. One of the simplest physically based cloth models over the last decade, and thus, the most likely to achieve real-time performances, is the mass-spring system [3].

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Fig. 1. A still from the film Sea Tails, showing Matisse' underwater kitemails.

46 In the mass-spring system, the deformable body is approx-  
 47 imated by a set of masses linked by springs in a  
 48 fixed topology. It is easy to implement, highly paral-  
 49 lelizable, and involves few computations. In addition  
 50 to this, there is a plethora of existing techniques for  
 51 non-real-time requirements, which the reader is en-  
 52 couraged to peruse [4].

53 All the above-mentioned approaches, however, suf-  
 54 fer from the same problem—to ensure stability, the  
 55 simulation has to be performed in very small time  
 56 steps making them very computationally intensive.  
 57 Various ways to overcome this problem have been  
 58 suggested. One model is the recent development of  
 59 neuro-animators [5], where after a learning period,  
 60 a large neural network can emulate a simple physi-  
 61 cal system. This recent approach has not been proven  
 62 practical for large coupled systems such as cloth. The  
 63 use of implicit integration, which can stably take large  
 64 time steps, has been proposed [6] in the context of  
 65 cloth animation. More recently, implicit approaches  
 66 to mass-spring systems are proposed by Meyer et al.  
 67 [7] in the context of VR environments. Although the  
 68 implicit method offers significantly reduced computa-  
 69 tional times, the kind of approximation does not ren-  
 70 der very visually accurate simulations.

71 To solve this computationally intensive problem,  
 72 we propose a distributed approach using the Grid

with its geographically dispersed processors linked by  
 high-speed interconnects. The rest of the paper will  
 proceed as follows. In Section 2, we will review our  
 physically based model, explaining the mass-spring  
 setup and the dynamics. Section 3 will discuss the  
 standalone implementation on a single machine quan-  
 titatively presenting the limitations and results. We  
 will extend this in Section 4 to a distributed networked  
 architecture, describing the details of the system and  
 an evaluation of results in comparison with the stan-  
 dalone version. Finally, in Section 5, we conclude  
 with some discussion on possible future work.

## 2. The physically based model

The physically based mass-spring cloth model [8,9]  
 is used to simulate the behavior of the kites. Each kite  
 is modeled as a cloth object that is approximated to  
 an array of masses and springs. Using the fundamen-  
 tal laws of dynamics, various forces acting on these  
 masses and springs are evaluated. These forces create  
 the movement of the individual masses in the network  
 and thus the deformation of the cloth as a whole is  
 simulated.

### 2.1. Mesh model

The kite model is composed of masses and springs,  
 where each kite is a grid of masses that form the  
 control points for the motion. Each mass-point in this  
 grid, say  $P_{ij}$  is connected to neighboring points with  
 springs. When two points get pulled further apart,  
 they experience a force pulling them together and  
 vice versa. There are three different kinds of springs:  
*structural springs* connect adjacent horizontal and  
 vertical points thus defining the rough structure of  
 the kite. These springs handle the compression and  
 traction stresses. However, these simple spring con-  
 nections alone are not enough to force the grid to  
 hold its shape causing it to shear. To counteract  
 this shearing, we add *shear springs* that also con-  
 nect these points diagonally. Fig. 2a shows the mesh  
 with the structural and shear springs. However, this  
 model is still incomplete. Once moving, the kites  
 tend to fray very easily as there is nothing to keep  
 the model from folding along the edges. So, we need  
 to add *bend springs* that resist folding and bend-

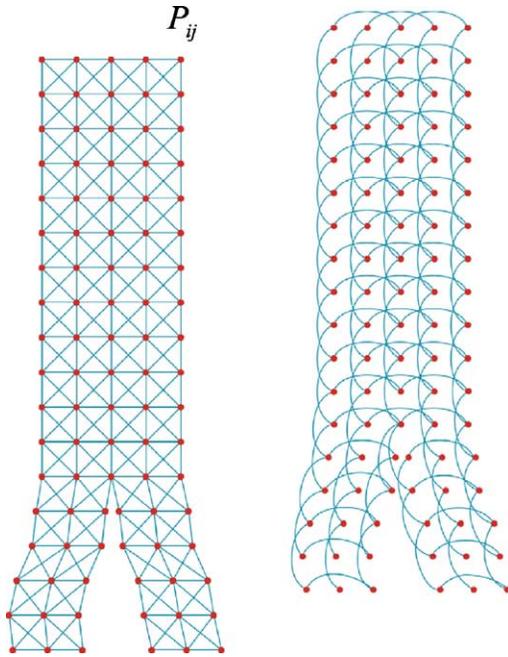


Fig. 2. The kite mesh model: (a) shear and structural springs, (b) bend springs.

116 ing. These bend springs connect every other point  
117 in the horizontal and vertical directions as shown in  
118 Fig. 2b.

119 The final mesh is a superimposition of the two  
120 meshes shown in Fig. 2.

## 121 2.2. Dynamics

122 The movement of kitetails is in turn determined by  
123 the dynamics of the mass-points or particles in the  
124 mesh. So, we want to individually model the motion  
125 of each particle  $P_{ij}$ . The problem is formulated as fol-  
126 lows: assuming we know the position  $x_{ij}$ , velocity  $v_{ij}$   
127 and resultant force  $F_{ij}$  acting on a particle  $P_{ij}$  at time  
128  $t$ , we want to compute the new position  $x_{ij}^{t+\Delta t}$ , of that  
129 particle after a small amount of time,  $\Delta t$ , has elapsed.

130 With the value of the force,  $F_{ij}$  and the mass  $m$  of  
131 the particle, we can obtain the acceleration  $a_{ij}$  of the  
132 particle as stated in the familiar Newtonian notation,  
133  $F = ma$ . So, at time  $t + \Delta t$ , the acceleration,  $a_{ij}^{t+\Delta t}$   
134 of the particle  $P_{ij}$  is given by

$$135 \quad a_{ij}^{t+\Delta t} = \frac{1}{m} F_{ij}^{t+\Delta t}$$

Integrating this acceleration with respect to time, we  
get the new velocity of the particle

$$v_{ij}^{t+\Delta t} = v_{ij}^t + \Delta t a_{ij}^{t+\Delta t}$$

Integrating again, we end up with the new position,  $x_{ij}$

$$x_{ij}^{t+\Delta t} = x_{ij}^t + \Delta t v_{ij}^{t+\Delta t}$$

The integration steps above give rise to the classical  
problem in simulation systems, that of *numerical instabil-*  
*ity*. This arises because the integration of a con-  
tinuous function is approximated by a discrete numeri-  
cal integrator. When the error between the approxi-  
mation and the real value gets too large, the numeri-  
cal simulation can fail. The choice of the integrator  
and the time step,  $\Delta t$ , are thus crucial. The first-order  
Euler's method [10] although very simple is subject  
to numerical instability. We instead use the midpoint  
method [10] which is more robust albeit computationally  
expensive. However, even with a robust integra-  
tor, there will be times when the simulation will be  
in danger of diverging or “blowing up”. The solution  
therefore is to take small time steps. However, smaller  
time steps mean more processing. If the processor is  
not fast enough, the simulation will be too slow to ob-  
serve anything visually interesting.

Now all that is left is to determine  $F_{ij}$  for a particle,  
which is the sum of all the different types of forces  
acting on it. There are two types of forces we model—  
*Internal* and *External* forces.

## 2.3. Internal forces

These are the forces determined by the physical  
properties of the cloth, i.e. the structural, shearing and  
bending forces and are a result of the tensions of these  
springs linking neighboring points of the cloth.

As a spring is compressed or expanded, it creates a  
force that is opposed to direction of the applied force.  
Assume we have two particles  $P_{ij}$  and  $P_{kl}$  connected  
by a spring,  $S$  which can be of any type (bend, shear  
or structural) as shown in Fig. 3.

The force contribution from spring  $S$ ,  $F_S$  is mathe-  
matically equated by the formula

$$F_S = k(L^t - L^0)(P_{ij} - P_{kl})$$

where  $k$  is the elasticity coefficient (set to 1000 N/m),  
 $L^t$  the length of the spring at time  $t$ ,  $L^0$  the rest length  
of the spring, i.e. at time 0.

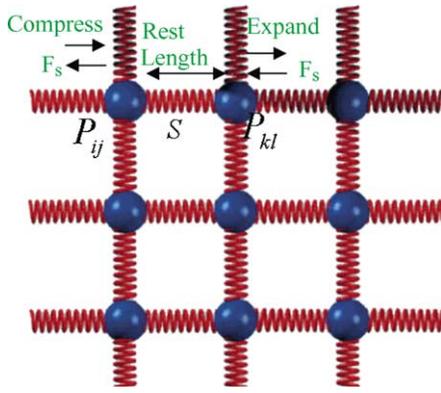


Fig. 3. Modeling the internal forces contributed by the springs.

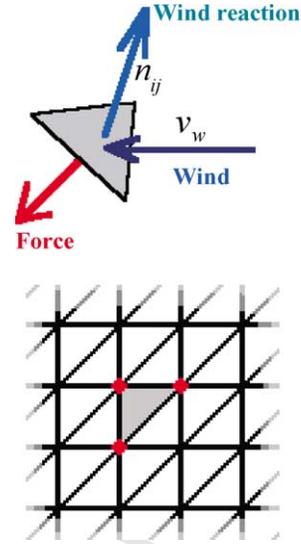


Fig. 4. Modeling wind force. The kite mesh is treated as a surface in order to compute the wind reaction.

179 This force  $F_s$  is accumulated for both  $P_{ij}$  and  $P_{kl}$  but  
 180 in opposing directions. This force computation process  
 181 is repeated for every spring linked to the point  $P_{ij}$ ,  
 182 and summed to give the final internal force vector  
 183  $F_{inij}$ .

#### 184 2.4. External forces

185 External forces are forces applied to the entire sys-  
 186 tem, which in our case are *gravity*, *wind resistance*  
 187 and *viscous drag*.

188 The kites suspended in the air will fall due to the  
 189 force of gravity pushing downwards to the earth. *Grav-*  
 190 *itational force*,  $F_{gr}$  acts on all particles and is modeled  
 191 as

$$192 F_{gr_{ij}} = mg$$

193 where  $m$  is the mass of the particle  $P_{ij}$ ,  $g$  the acceler-  
 194 ation of gravity, taken to be  $9.8 \text{ m/s}^2$ .

195 Since this force depends on mass, a kite with a  
 196 larger mass will fall faster than one with a smaller  
 197 mass.

198 *Wind force*,  $F_w$  acts on a particle depending on  
 199 its surface normal. As shown in Fig. 4, this force is  
 200 greatest when the surface of the kite and the wind  
 201 vector are perpendicular since this gives the great-  
 202 est cross-sectional area and therefore greater air resis-  
 203 tance.

204 The wind force is calculated as follows:

$$205 F_{w_{ij}} = \mu_w n_{ij} (v_w - v_{ij}) n_{ij}$$

where  $\mu_w$  is a user-specified viscosity constant for the  
 wind,  $v_w$  the wind vector,  $n_{ij}$  the normal to the surface  
 at point  $P_{ij}$ .

In order to compute the surface normals, the kite  
 mesh is tessellated to form triangles. The wind force  
 is calculated on each of these triangles individually.  
 At each point,  $P_{ij}$  of the cloth, the sum of the effect  
 of the wind on the surrounding triangles is calculated.  
 Since the mesh is constantly changing, this normal  
 computation needs to be done every frame.

The *viscous drag*,  $F_d$  is a damping force applied to  
 each particle and is directly proportional and opposite  
 to the velocity of the moving particle. We use this  
 drag to model the loss of mechanical energy of the  
 cloth and also add numerical stability to the system,  
 ensuring that the particles will not bounce around too  
 much:

$$223 F_{d_{ij}} = -\mu_d v_{ij}$$

where  $\mu_d$  is a user-specified damping coefficient (set  
 to be  $20.0 \text{ N/m}$ ).

The internal and external force vectors are summed  
 to give the final force acting on the point  $P_{ij}$

$$228 F_{ij} = F_{inij} + F_{gr_{ij}} + F_{w_{ij}} + F_{d_{ij}}$$

232 It is the combination of the above forces that gives the  
233 kite its smooth, fluttering motion.

### 231 3. Standalone implementation

232 The preliminary standalone version of “Kites fly-  
233 ing In and Out of Space” was demonstrated on the  
234 CAVE system at Virginia Tech as part of the *Moun-  
235 tain Lake Workshop* [11] which is a collaborative,  
236 community-based art project drawing on the customs,  
237 environmental resources, and technology of the New  
238 River Valley and the Appalachian region, USA.

#### 239 3.1. CAVE virtual reality

240 The CAVE (CAVE Automatic Virtual Environment)  
241 is a projection-based virtual reality system [12]. In  
242 contrast to head-mounted display VR systems, where  
243 the user views a virtual world through small video  
244 screens attached to a helmet, in projection-based VR  
245 large, fixed screens are used to provide a panoramic  
246 display without encumbering the user. The CAVE  
247 is a 10 ft cubed room. Stereoscopic images are  
248 rear-projected onto the walls creating the illusion  
249 that 3D objects exist with the user in the room. The  
250 user wears liquid crystal shutter glasses to resolve the  
251 stereoscopic imagery. An electromagnetic tracking  
252 sensor attached to the glasses allows the CAVE system  
253 to determine the location and orientation of the user’s  
254 head. This information is used by the Silicon Graph-  
255 ics Onyx that drives the CAVE to render the imagery  
256 from the user’s point of view. The user can physically  
257 walk around an object that appears to exist in 3D in  
258 the middle of the CAVE. The user holds a wand that  
259 is also tracked and has a joystick and three buttons  
260 for interaction with the virtual environment. Typically  
261 the joystick is used to navigate through environments  
262 that are larger than the CAVE itself. The buttons  
263 can be used to change modes, or bring up menus in  
264 the CAVE, or to grab virtual objects. Speakers are  
265 mounted to the top corners of the CAVE structure to  
266 provide sounds from the virtual environment.

267 Software support for the CAVE comes in the form  
268 of the CAVE library. Applications are built on top of  
269 the CAVE library, which controls the display, tracking,  
270 and input devices. OpenGL which is an industry stan-  
271 dard Graphics API was used to render these kitetails.

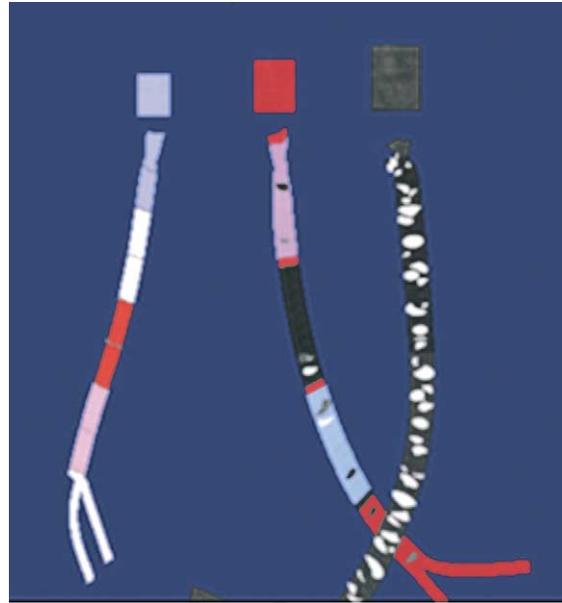


Fig. 5. A snapshot of the standalone kite application.

#### 3.2. Interacting with the kitetails

272

273 Fig. 5 shows a snapshot of the application that has  
274 three kites in the scene. The imagery for these tails  
275 was scanned from the physical kitetails and texture  
276 mapped onto the kitetail mesh. These kites are an-  
277 chored by their two end points and start off horizontal  
278 thus free falling under the influence of gravity. Using  
279 the CAVE wand, the user can grab on to a kite head to  
280 move or change its imagery. At any point, the move-  
281 ment of these tails can be dynamically controlled by  
282 injecting a global wind into the system using the wand.  
283 The wind strength is constant in space and its direction  
284 is determined by the wand orientation. The structural  
285 attributes of a kite like the stiffness, length, width and  
286 the visual attributes like its texture maps can be spec-  
287 ified at run-time by the user using a configuration file.  
288 Each kite has a dimension of 2 ft × 30 ft in the virtual  
289 space and is modeled by about 250 mass-points.

#### 3.3. Results and findings

290

291 The number of iterations computed by the simu-  
292 lator per second or the simulation rate can be used  
293 as a metric to determine its performance. A higher

294 value indicates a faster simulation. The simulation rate  
 295 for one kite in our application was 125 iterations per  
 296 second and this varied inversely with the number of  
 297 kites in the system. The compute time for each iteration  
 298 was therefore, 8 ms. The time step was chosen  
 299 to be 5 ms, to be close enough to the actual iteration  
 300 time. A smaller time step simulation would have  
 301 been more stable but also more compute intensive.  
 302 We instead decided to “fly” three kites. Although this  
 303 slows the simulation rate for each kite to 41.7 iterations  
 304 per second, the results were visually more interesting.  
 305 The simulation ran on a SGI ONYX Infinite  
 306 Reality with eight 195 MHz MIPS R10000 processors  
 307 and 2048 MB main memory size.

#### 308 4. Distributed simulation

309 As seen in the previous section, the small time-step  
 310 requirement calls for additional computing power. The  
 311 GRID with its distributed computing resources connected  
 312 by high-speed networks proves to be an elegant and  
 313 scalable solution to circumvent the processing constraints  
 314 imposed by a single machine. Although, theoretically,  
 315 any speed network can be used

316 as long as it supports the underlying network protocols.  
 317

318 Fig. 6 pictorially describes the structure of the distributed  
 319 system. We decouple the simulation from the display  
 320 procedures. This results in many simulation nodes  
 321 all across the GRID and one display node, in our case  
 322 the CAVE machine that displays the results of the  
 323 simulation. These distributed software components  
 324 communicate using a middleware QUANTA [13] that  
 325 was developed within EVL. QUANTA is a rich collection  
 326 of network programming tools for optimizing data  
 327 sharing over high-speed networks. Following is a  
 328 brief description of the distributed components used:

- *kiteServer* is an implementation of the QUANTA  
 329 database module that provides a simple two-field  
 330 database, associating arbitrary chunks of binary data  
 331 with character string keys. The keys are treated like  
 332 Unix directory paths, so that a hierarchical arrangement  
 333 of data is possible. When a client connects to a  
 334 QUANTA database, it can make asynchronous requests  
 335 to fetch particular keys’ values, and it can store  
 336 new values for keys. Stored data is automatically  
 337 reflected to all other clients by the database server.  
 338 In our application, the shared data is the wind

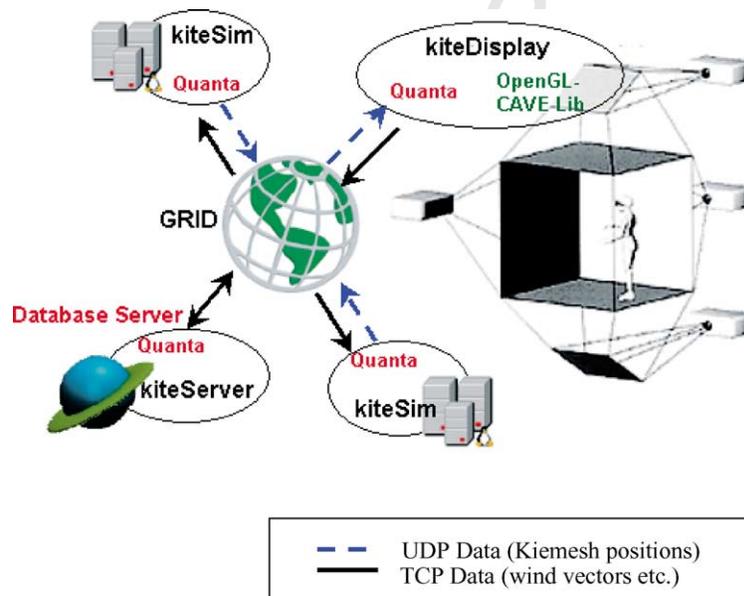


Fig. 6. The distributed model for the kite application. Typically, there are many simulation nodes, one display client (e.g. CAVE) and a server that maintains state information about user-interaction. These components are connected across the GRID or any high-speed network.

339 direction specified as a 3-float array. This wind direction  
340 direction vector is received from the CAVE display  
341 client as a result of any user-interaction and broad-  
342 cast to all the *kiteSim* nodes to update their simula-  
343 tion.

344 • *kiteSim* is the simulation server that typically runs  
345 on a linux machine. Each machine's simulation  
346 loop is dedicated to computing the positions for  
347 one kitetail. These computed positions are directly  
348 transmitted through a UDP socket to the display  
349 client running in the CAVE. The network loop,  
350 in turn, uses the database client to listen to any  
351 changes in the wind direction broadcast from the  
352 *kiteServer*.

353 • *kiteDisplay* is the display client that displays the  
354 kitetails and handles any user-interaction. The kite  
355 position data for each kite mesh is read from the dis-  
356 tributed servers using QUANTA UDP and displayed  
357 texture mapped with the images. In addition, any  
358 user-interaction events to move the kites or change  
359 the imagery are handled here. When the user injects

360 wind into the system, the wind direction vector is  
361 sent to the *kiteServer* using a QUANTA database  
362 client, which is then broadcasted to all the simula-  
363 tion nodes. The OpenGL/CAVElib is used for the  
364 user-interaction and display.

#### 4.1. Implementation 365

366 The network testbed used in IGRID2002 [14] is  
367 shown in Fig. 7.

368 The CAVE in SARA, Amsterdam was our display  
369 client with simulation servers distributed across the  
370 globe in Chicago, Canada, Japan, and Virginia  
371 to stream the kitetails positions. The test itself involved  
372 a gamut of networking infrastructures ranging from  
373 the 10 Gbps optical link between Starlight/Surfnet to  
374 the more commonplace Internet2. The network band-  
375 width indicated for each machine is the bandwidth  
376 required if the kite position data were streamed to the  
377 network at the simulation rate (which is the maximum  
378 possible rate).

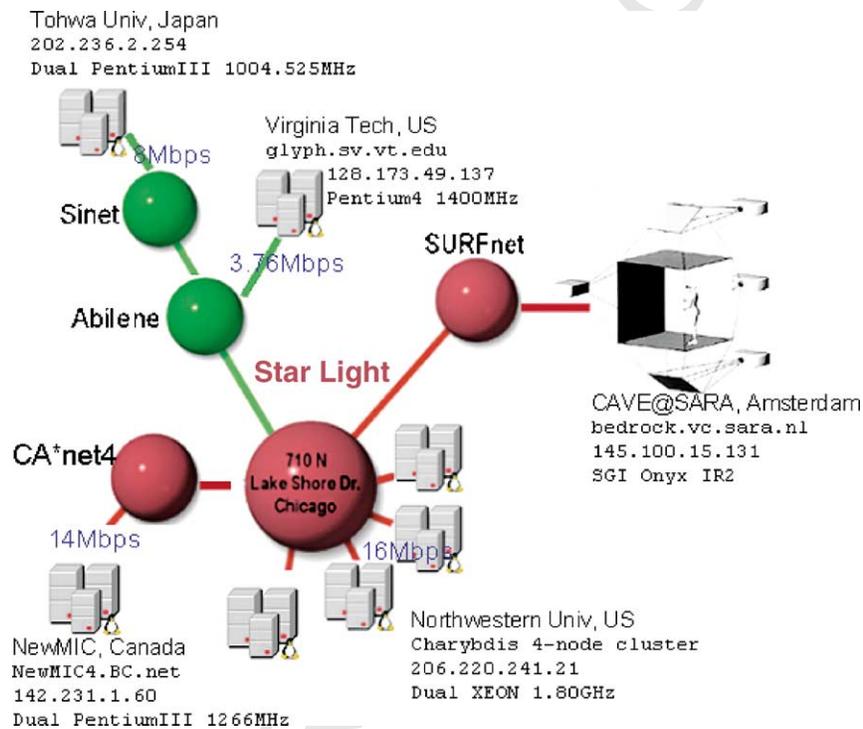


Fig. 7. The network testbed used for IGRID 2002. The display client was the CAVE at SARA, Amsterdam with the simulation running across the GRID in Chicago, Canada, Japan and Virginia.

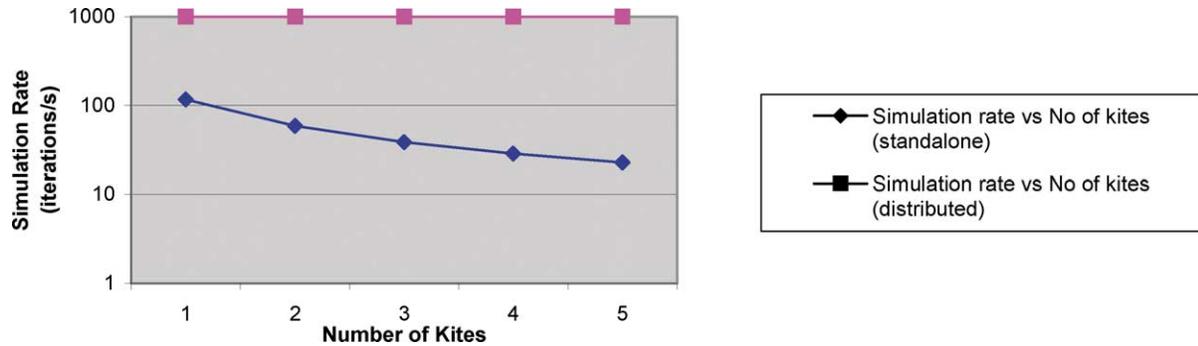


Fig. 8. Results showing that the simulation rate is significantly higher and is independent of the number of kites for the distributed version. The scale of the simulation rate is logarithmic.

#### 379 4.2. Results

380 The graph Fig. 8 shows in logarithmic scale the  
 381 performance of the distributed system vis-à-vis its  
 382 standalone version. Two conclusions follow from this  
 383 graph. Firstly, since the simulations were running on  
 384 high-end PCs, the simulation rate for a kite in the dis-  
 385 tributed version (1000 iterations per second) is signifi-  
 386 cantly higher than that of the standalone (125 itera-  
 387 tions per second). The compute time for one iteration  
 388 is 1 ms. Hence, smaller time steps could be used which  
 389 gave us a stable simulation. Since the time step is equal  
 390 to the compute time for an iteration, the movement of  
 391 the kites is also more natural and cloth-like. Secondly,  
 392 in the distributed version, the simulation rate is inde-  
 393 pendent of the number of kites. Of course, the limiting  
 394 factor in this case could be the network bandwidth.  
 395 The network bandwidth used by each simulation was  
 396 about 1.1 Mbits/s. Although the simulation ran as fast  
 397 as the processor (1000 iterations per second), the re-  
 398 sults were only streamed to the network at 40 frames  
 399 per second which suffices for the graphics update.

400 Fig. 9 shows a screenshot of the distributed version  
 401 that was demonstrated in the CAVE at SARA, Am-  
 402 sterdam for IGRID2002. There are 12 kites in total,  
 403 “flying” from all the remote simulation sites. The kites  
 404 crisscross each other in the virtual space since there are  
 405 no spatial limits imposed on them. As the movement  
 406 of the kites is dependant on the network, they become  
 407 visual metaphors for their underlying network perfor-  
 408 mance and parameters. A slow moving kite for in-  
 409 stance, signifies a low-bandwidth link and vice versa.

410 The response time of the kites to wind interaction in-  
 411 dicates their network latency. Since the kitetail data  
 412 is split into segments and transmitted across as UDP  
 413 packets, any visual jaggedness in the reconstruction at  
 414 the display side represents the segment size and the  
 415 UDP packet loss.



Fig. 9. A screenshot of the distributed version of the kites application. The yellow particle traces represent the wind forces added.

## 416 5. Conclusion and future work

417 The application of Grid computing for the Arts and  
 418 Humanities is still in its infancy. It is our hope that this  
 419 first piece of GRID Art will encourage others to think  
 420 of creative new ideas for exploiting GRID technolo-  
 421 gies. As part of our future work, we would like to make  
 422 the application collaborative so that CAVEs around  
 423 the world could potentially view this application. For  
 424 complex physical interactions like collision detection,  
 425 we would like a synchronous mode that would involve  
 426 more communication between the simulation nodes.  
 427 In order to further optimize, we could use some im-  
 428 plicit methods of integration for the simulation. We  
 429 would also like to improve the rendering quality to vi-  
 430 sually model the interaction of kitetails with different  
 431 media as this is the focus of Matisse' works.

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